



Suppose  $f(t)$  is square integrable  $f \in L^2([0,1])$

i.e.  $\int_0^1 |f(t)|^2 dt < \infty$  (finite energy)

If  $f(t)$  periodic, period 1,  $f \in L^2([0,1])$

$$\hat{f}(k) = \int_0^1 f(t) e^{2\pi i k t} dt$$

$$\int_0^1 \left| f(t) - \sum_{k=-\infty}^{n-1} \hat{f}(k) e^{2\pi i k t} \right|^2 dt \rightarrow 0 \text{ as } n \rightarrow \infty$$

only get these convergence results with generalization of Integral  
what I learn is Riemann Integral

$$\int_0^1 e^{2\pi i m t} \cdot e^{2\pi i n t} dt = \begin{cases} 0 & m \neq n \\ 1 & m = n \end{cases}$$

Consequence for introducing "geometry" into  $L^2([0,1])$   
Allow one to define "orthogonality" via inner product

Assume  $f, g$  are square integrable, inner product  $\langle f, g \rangle = \int_0^1 f(t) \bar{g}(t) dt$

$$\text{Norm } \|f\|^2 = \int_0^1 (f(t))^2 dt$$

$f, g$  are orthogonal if  $\langle f, g \rangle = 0$

Pythagorean theorem:

$$f \perp g \iff \|f+g\|^2 = \|f\|^2 + \|g\|^2 \iff \langle f, g \rangle = 0$$

Vector  $\rightarrow$  function

Reasoning by strong analogy

$$\langle e^{2\pi i m t}, e^{2\pi i n t} \rangle = \begin{cases} 0 & m \neq n \\ 1 & m = n \end{cases} \quad e^{2\pi i m t}'s \text{ are orthonormal}$$

Fourier coefficients are projections of  $f(t)$  on  $e^{2\pi i m t}$

Vector

inner product

$$\langle u, v \rangle = \sum_i u_i v_i$$

norm

$$\|u\|_2^2 = \sum_i |u_i|^2$$

orthogonality

$$U \perp V \iff \langle U, V \rangle = 0$$

projection

$$\underbrace{\frac{v}{\|v\|}}_{u} \cdot \underbrace{\langle u, v \rangle}_{\langle u, v \rangle \cdot u}$$

function

$$\langle f, g \rangle = \int_0^1 f(t) \bar{g}(t) dt$$

$$\|f\|^2 = \int_0^1 |f(t)|^2 dt$$

$$f \perp g \iff \langle f, g \rangle = 0$$

$$\langle f, g \rangle \neq 0$$

$$f(t) = \sum_{k=-\infty}^{+\infty} \hat{f}(k) \cdot e^{2\pi i k t} = \sum_{k=-\infty}^{+\infty} \langle f, e^{2\pi i k t} \rangle \cdot e^{2\pi i k t}$$

$\{e^{2\pi i k t} \mid -\infty < k < +\infty\}$  form orthonormal basis for square integrable periodic functions

Kaleys Identity:

$$\int_0^t |f(\epsilon)|^2 d\epsilon = \sum_{k=-\infty}^{+\infty} |\hat{f}(k)|^2$$

norm of a vector = sum of square of components

$$f(t) = \sum_{k=-\infty}^{+\infty} \hat{f}(k) e^{2\pi i k t}$$

↓

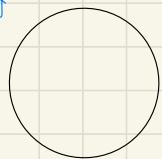
$f(t)$  has energy  $\int_0^t |f(\epsilon)|^2 d\epsilon$       each component has energy  $|\hat{f}(k)|$  ( $e^{2\pi i k t}$  has energy 1)

Application to heat flow

have region in space with initial temperature distribution  $f(x)$  at  $t=0$  ( $x$  is space variable)

how does temperature change in position and time

e.g. heated ring



given initial temperature function

$u(x,t)$  is the temperature at position  $x$  at time  $t$

periodicity: symmetry in space  $\begin{cases} f(x+t) = f(x) \\ u(x+1, t) = u(x, t) \end{cases}$

$$u(x,t) = \sum_{k=-\infty}^{+\infty} C_k(t) e^{2\pi i k x} : \text{what are } C_k(t)$$

$$\begin{cases} \frac{du}{dt} = a \cdot \frac{d^2 u}{dx^2} & (\text{heat equation / diffusion equation}) \\ = \frac{1}{2} \frac{d^2 u}{dx^2} & \text{choose } a = \frac{1}{2} \end{cases}$$

$$\sum_{k=-\infty}^{+\infty} \frac{dC_k}{dt} e^{2\pi i k x} = \sum_{k=-\infty}^{+\infty} C_k(t) \cdot (2\pi i k)^2 e^{2\pi i k x}$$

|| equate coefficients (each function has 1 F.S.)

$$\frac{dC_k}{dt} = C_k(t) \cdot (2\pi i k)^2$$

$$C_k(t) = C_k(0) \cdot e^{(2\pi i k)^2 t}$$